

# Dynamical Systems and Quasiconformal Mappings:

A Course Given in Department of Mathematics at CUNY Graduate Center  
Spring Semester of 2009, Friday, 9:30am-11:30am

Yunping Jiang

## Lecture VI, March 6, 2009

### Horseshoe Maps

Let  $R$  be a rectangle in  $\mathbb{R}^2$  with two horizontal sides and two vertical sides. Suppose  $f : R \rightarrow f(R) \subset \mathbb{R}^2$  is a diffeomorphism such that  $R \cap f(R)$  are two separate rectangles  $R_0$  and  $R_1$  such that their horizontal sides parallel to the horizontal sides of  $R$  and their vertical sides are parts of the vertical sides of  $R$ . Let  $R^0$  and  $R^1$  be two components of  $f^{-1}(R) \cap R$ . Suppose  $R^0$  and  $R^1$  are also two rectangles whose vertical sides are parallel to those vertical sides of  $R$  and whose horizontal sides are parts of the horizontal sides of  $R$ . Suppose  $f : R^0 \rightarrow R_0$  and  $f : R^1 \rightarrow R_1$  are hyperbolic affine maps, contracting in the vertical direction and expanding in the horizontal direction. This map is called a Smale horseshoe map. Let  $\Lambda$  be the maximal invariant set of  $f$  in  $R$ . Then

$$\Lambda = \bigcap_{n=-\infty}^{\infty} f^n(R).$$

The intersection

$$R \cap f(R) \cap f^2(R) = (R_0 \cap R_1) \cap f^2(R)$$

consists of four rectangles  $R_{ij}$  for  $i, j \in \{0, 1\}$  whose horizontal sides parallel to the horizontal sides of  $R$  and their vertical sides are parts of the vertical sides of  $R$ . Inductively,  $\bigcap_{k=0}^n f^k(R)$  consists of  $2^n$  thin rectangles whose horizontal sides parallel to the horizontal sides of  $R$  and their vertical sides are parts of the vertical sides of  $R$ . Let us denote these rectangles as  $R_{i_0 i_1 \dots i_{n-1}}$  for  $i_k \in \{0, 1\}$ ,  $k = 0, 1, \dots, n-1$ . Here we label them as  $f(R_{i_0 \dots i_{n-1}}) = R_{i_1 \dots i_{n-1}}$ . Then we have that

$$R_{i_0 \dots i_{n-1}} \subset R_{i_0 \dots i_{n-2}} \subset \dots \subset R_{i_0}.$$

This implies that

$$R_{i_0 \dots i_{n-1} \dots} = \bigcap_{k=0}^{\infty} R_{i_0 \dots i_{n-1}}$$

is a horizontal line in  $R$ . Then

$$R_+ = \bigcap_{k=0}^{\infty} f^k(R)$$

consists of uncountably many horizontal lines.

Similarly, the intersection

$$R \cap f^{-1}(R) \cap f^{-2}(R) = (R^0 \cap R^1) \cap f^{-2}(R)$$

consists of four rectangles  $R^{ij}$  for  $i, j \in \{0, 1\}$  whose vertical sides parallel to the vertical sides of  $R$  and and their horizontal sides are parts of the horizontal sides of  $R$ . Inductively,  $\bigcap_{k=1}^n f^{-k}(R)$  consists of  $2^n$  thin rectangles whose horizontal sides parallel to the horizontal sides of  $R$  and and their vertical sides are parts of the vertical sides of  $R$ . Let us denote these rectangles as  $R^{j_{n-1}j_{n-2}\cdots j_1}$  for  $j_k \in \{0, 1\}$ ,  $k = 1, \dots, n-1$ . Here we label them as  $f^{-1}(R_{j_{n-1}\cdots j_2 j_1}) = R_{j_{n-1}\cdots j_2}$ . Then we have that

$$R_{j_{n-1}j_{n-2}\cdots j_1} \subset R_{j_{n-2}} \subset \cdots \subset R_{j_1}.$$

This implies that

$$R_{\dots j_{n-1}\cdots j_1} = \bigcap_{k=1}^{\infty} R_{j_{n-1}\cdots j_1}$$

is a vertical line in  $R$ .

$$R_- = \bigcap_{k=0}^{\infty} f^k(R)$$

consists of uncountably many horizontal lines.

Give any  $w_+ = i_0 \cdots i_{n-1} \cdots$ ,  $R_{w_+} \cap R_-$  is a Cantor set on the line  $R_{w_+}$ . Give any  $w_- = \cdots j_{n-1} \cdots j_1$ ,  $R_{w_-} \cap R_+$  is a Cantor set on the line  $R_{w_-}$ . Then

$$\Lambda = R_+ \cap R_-$$

can be thought as a product of two Cantor set on a horizontal line and on a vertical line.

**Theorem 1.** *Consider the dynamical system  $f : \Lambda \rightarrow \Lambda$ . Then the set of periodic points of  $f$  are dense in  $\Lambda$ . And  $Per_n(f)$  contains  $2^n$  different points. Furthermore,  $f : \Lambda \rightarrow \Lambda$  is topologically mixing, that is, for any two open sets  $U$  and  $V \subset \Lambda$ , there exists an integer  $N = N(U, V) > 0$  such that  $f^n(U) \cap V$  for is nonempty for every  $n > N$ .*

### One-side full shift on symbolic space.

In one-dimensional expanding maps, we used symbols to label intervals and points. Here we study a symbolic dynamical system. Through the study of symbolic dynamical systems, we can have a better understanding of dynamical systems.

Let  $S = \{0, 1, \dots, d-1\}$  be a set with the discrete topology. Consider the space

$$\Sigma = \Sigma_d = \prod_{k=0}^{\infty} S = \{w = i_0 i_1 \cdots i_{k-1} \cdots \mid i_k \in S, k = 0, 1, \dots\}.$$

We give  $\Sigma$  the product topology as follows. Give any

$$w = i_0 i_1 \cdots i_{k-1} \cdots .$$

Define

$$[w]_n = \{w' = i_0 i_1 \cdots i_{n-1} j_n \cdots j_{k-1} \cdots \in \Sigma\}.$$

It is called an  $n$ -cylinder of  $\Sigma$ . We call a cylinder an open set. The set of all cylinders forms a topological basis. Then  $\Sigma$  with this topology is a topological space.

We can also introduce a metric on  $\Sigma$ :

$$d(w, w') = \sum_{k=1}^{\infty} \frac{|i_{k-1} - j_{k-1}|}{d^k}$$

if  $w = i_0 \cdots i_k \cdots$  and  $w' = j_0 \cdots j_{k-1} \cdots$ . One can check that  $d(\cdot, \cdot)$  is a metric on  $\Sigma$  and the topology induced from this metric on  $\Sigma$  is the same as the product topology.

One side full shift is defined as

$$\sigma : \Sigma \rightarrow \Sigma; \quad \sigma(i_0 i_1 \cdots i_{k-1} \cdots) = i_1 \cdots i_{k-1} \cdots .$$

Then it is a continuous map. The dynamical system  $\sigma$  is called a symbolic dynamical system.

Suppose  $I = [0, 1]$  and  $I_0 = [0, b_0]$ ,  $I_1 = [a_1, b_1]$ ,  $\dots$ ,  $I_{d-1} = [a_{d-1}, b_{d-1}]$  where

$$a_0 = 0 < b_0 < a_1 < b_1 < \cdots < a_{d-1} < b_{d-1}.$$

Let

$$f : U = \cup_{i=0}^{d-1} I_i \rightarrow I$$

be a map such that  $f : I_i \rightarrow I$  is a  $C^1$ -diffeomorphism for each  $i = 0, \dots, d-1$ . We say  $f$  is expanding if there are two constants  $C > 0$  and  $\lambda > 1$  such that

$$|(f^n)'(x)| \geq C\lambda^n$$

for all  $x \in I$  and  $n > 0$  such that  $f^{n-1}(x)$  is defined. Let

$$\Lambda = \cap_{k=0}^{\infty} f^{-k}(I)$$

be the maximal invariant set of  $f$ .

**Theorem 2.** *Suppose  $f : U \rightarrow I$  is a  $C^1$  expanding map. Then  $f : \Lambda \rightarrow \Lambda$  is topologically conjugate to the symbolic dynamical system  $\sigma : \Sigma \rightarrow \Sigma$ . That is, there is a homeomorphism  $h : \Sigma \rightarrow \Lambda$  such that*

$$f \circ h = h \circ \sigma$$

on  $\Sigma$ .

*Proof.* For each  $0 \leq i \leq d-1$ , let  $g_i = (f|_{I_i})^{-1} : I \rightarrow I_i$ . For each  $w_m = i_0 i_1 \cdots i_{m-1}$ , define

$$g_{w_m} = g_{i_0} \circ \cdots \circ g_{i_{m-1}}$$

and

$$I_{w_m} = g_{w_m}(I).$$

Then for every  $w = w_m \cdots \in \Sigma$ ,

$$\cdots \subseteq I_{w_m} \subset I_{w_{m-1}} \subset \cdots \subset I_{w_1}.$$

Since  $f$  is expanding,  $|I_{w_m}| \rightarrow 0$  as  $m \rightarrow \infty$ . So

$$I_w = \bigcap_{m=1}^{\infty} I_{w_m} = \{x_w\}$$

contains only one point. Define  $h : \Sigma \rightarrow \Lambda$  as  $h(w) = x_w$ , we have that  $h$  is a homeomorphism and

$$f \circ h = h \circ \sigma$$

on  $\Sigma$ . □

**Lemma 1.** *All periodic points of period  $m > 0$  are*

$$\text{Per}_m(f) = \{x_w \mid w = (i_0 \cdots i_{m-1})^\infty\}.$$

**Exercise 1.** *Prove that  $\sigma$  (as well as  $f$ ) is topologically transitive but not minimal.*

**Exercise 2.** *Find all proper closed subsets  $A$  contains only finitely many points in  $\Sigma$  which is forward invariant, that is,  $\sigma(A) \subseteq A$ . Find a proper closed subset  $A$  contains infinitely many points in  $\Sigma$  which is forward invariant.*

## Gibbs Measures

If a physical system of  $n$  states with the energies of these states are  $E_1, \dots, E_n$ . Suppose that this system is put in contact with a much larger “heat source” which is at temperature  $T$ . Energy is therefore allowed to pass between the original system and the heat source. Suppose the temperature  $T$  of the heat source remains constant. It is a

physical fact derived in statistical mechanics that the probability  $p_j$  that state  $j$  occurs is given by the Gibbs distribution

$$p_j = \frac{e^{-\beta E_j}}{\sum_{i=1}^n e^{-\beta E_i}}$$

where  $\beta = \frac{1}{kT}$  and  $k$  is a physical constant. This is the starting point for the thermodynamical formalism. However, thermodynamical formalism is a mathematical subject which studies Gibbs measures for more general systems.

Consider the symbolic dynamical system  $\sigma : \Sigma \rightarrow \Sigma$  (or  $f : \Lambda \rightarrow \Lambda$ ). A probability measure  $\mu$  on  $\Sigma$  is called an invariant measure if

$$\mu(A) = \mu(\sigma^{-1}(A))$$

for any Borel set  $A$ . A function  $\phi : \Sigma \rightarrow \mathbb{R}$  is called positive if  $\phi > 0$  and Hölder continuous if there are two constant  $C > 0$  and  $0 < \tau < 1$  such that

$$|\phi(w) - \phi(w')| \leq C\tau^m$$

for any  $w, w' \in \Sigma$  such that their first  $m$  digits are the same. For a positive and Hölder continuous function  $\phi$ , the function  $\psi = \log \phi$  is also Hölder continuous.

**Theorem 3.** *Consider the system  $(\sigma, \Sigma, \log \phi)$  where  $\phi$  is a positive Hölder continuous function. There is a unique probability measure  $\mu = \mu_{\log \phi}$  such that*

$$C^{-1} \leq \frac{\mu([w_m])}{\exp[-mP + \sum_{k=0}^{m-1} \log \phi(\sigma^k(w))]} \leq C.$$

where  $C > 0$  is a constant and where  $[w_m]$  is any cylinder of  $\Sigma$  and  $w \in [w_m]$ . Here  $P = P(\log \phi)$  is a constant depending on  $\log \phi$  and is called the pressure for the system.

**Remark 1.** *If we consider  $f : \Lambda \rightarrow \Lambda$  as a  $C^{1+\alpha}$  expanding map, let  $h : \Sigma \rightarrow \Lambda$  be the conjugacy from  $\Sigma \rightarrow \Lambda$ . Consider  $\phi(w) = 1/|f'(h(w))|$ . Then it is a positive Hölder continuous function on  $\Sigma$ . Consider*

$$t \log \phi(w) = -t \log |f'(h(w))|.$$

*Then the pressure  $P(-t \log |f'(h(w))|)$  is a function of  $t$ . There is a unique number  $0 < t_0 < 1$  such that  $P(-t \log |f'(h(w))|) = 0$ . This number is just the Hausdorff dimension of  $\Lambda$ . For more details about Gibbs theory, you can refer to my lecture notes (download site: <http://qcpages.qc.cuny.edu/~yjiang/HomePageYJ/Download/JiangNJLectureIFM.pdf>) "Nanjing Lecture Notes In Dynamical Systems. Part One: Transfer Operators in Thermodynamical Formalism."*